

Chapter 10

Local flows

We met local flows and integral curves in Chapter 6. Given a vector field v , write its local flow as ϕ_t .

- The collection ϕ_t for $t < \epsilon$ (for some $\epsilon > 0$, or alternatively for $t < 1$) is a **one-parameter group of local diffeomorphisms**. \square

Consider the vector field in a neighbourhood U of a point $Q \in \mathcal{M}$. Since $\phi_t : U \rightarrow \mathcal{M}, Q \mapsto \gamma_Q(t)$ is **local diffeomorphism**, i.e. diffeomorphism for sufficiently small values of t , we can use ϕ_t to push forward vector fields. At some point P we have the curve $\phi_t(P)$. We push forward a vector field at $t = \epsilon$ to $t = 0$ and compare with the vector field at $t = 0$.

We recall that for a map $\varphi : \mathcal{M}_1 \rightarrow \mathcal{M}_2$ the pullback of a function $f \in C^\infty(\mathcal{M}_2)$ is defined as

$$\varphi^* f = f \circ \varphi : \mathcal{M}_1 \rightarrow \mathbb{R}, \quad (10.1)$$

and $\varphi^* f \in C^\infty(\mathcal{M}_1)$ if φ is C^∞ .

The pushforward of a vector v_P is defined by

$$\varphi_* v_P(f) = v_P(f \circ \varphi) = v_P(\varphi^* f) \quad (10.2)$$

$$v_P \in T_P \mathcal{M}_1, \quad \varphi_* v_P \in T_{\varphi(P)} \mathcal{M}_2. \quad (10.3)$$

If φ is a diffeomorphism, we can define the pushforward of a vector field v by

$$\begin{aligned} \varphi_* v(f)|_{\varphi(P)} &= v(f \circ \varphi)|_P \\ \text{i.e.} \quad \varphi_* v(f)|_Q &= v(f \circ \varphi)|_{\varphi^{-1}Q} \\ &= v(\varphi^* f)|_{\varphi^{-1}Q}. \end{aligned} \quad (10.4)$$

We can rewrite this definition in several different ways,

$$\begin{aligned}(\varphi_*v)(f) &= v(f \circ \varphi) \circ \varphi^{-1} \\ &= (\varphi^{-1})^*(v(f \circ \varphi)) \\ &= (\varphi^{-1})^*(v(\varphi^*f)).\end{aligned}\tag{10.5}$$

- If $\varphi : \mathcal{M}_1 \rightarrow \mathcal{M}_2$ is not invertible, φ_*v is not a vector field on \mathcal{M}_2 . If φ^{-1} exists but is not differentiable, φ_*v is not differentiable. But there are some φ and some v such that φ_*v is a differentiable vector field, even if φ is not invertible or φ^{-1} is not differentiable. Then v and φ_*v are said to be **φ -related**. \square

Proposition: Given a diffeomorphism $\varphi : \mathcal{M}_1 \rightarrow \mathcal{M}_2$ (say both C^∞ manifolds) the pushforward φ_* is an isomorphism on the Lie algebra of vector fields, i.e.

$$\varphi_*[u, v] = [\varphi_*u, \varphi_*v].\tag{10.6}$$

Proof:

$$\begin{aligned}\varphi_*[u, v](f) &= [u, v](f \circ \varphi) \circ \varphi^{-1} \\ &= u(v(f \circ \varphi)) \circ \varphi^{-1} - u \leftrightarrow v,\end{aligned}\tag{10.7}$$

while

$$\begin{aligned}[\varphi_*u, \varphi_*v](f) &= \varphi_*u(\varphi_*v(f)) - u \leftrightarrow v \\ &= u(\varphi_*v(f) \circ \varphi) \circ \varphi^{-1} - u \leftrightarrow v \\ &= u((v(f \circ \varphi) \circ \varphi^{-1}) \circ \varphi) \circ \varphi^{-1} - u \leftrightarrow v \\ &= u(v(f \circ \varphi)) \circ \varphi^{-1} - u \leftrightarrow v.\end{aligned}\tag{10.8}$$

\square

- A vector field v is said to be **invariant** under a diffeomorphism $\varphi : \mathcal{M} \rightarrow \mathcal{M}$ if $\varphi_*v = v$, i.e. if $\varphi_*(v_P) = v_{\varphi(P)}$ for all $P \in \mathcal{M}$. \square

We can write for any $f \in C^\infty(\mathcal{M})$

$$\begin{aligned}(\varphi_*v)(f) &= (\varphi^{-1})^*(v(\varphi^*f)) \\ \Rightarrow \varphi^*((\varphi_*v)(f)) &= v(\varphi^*f), \\ \Rightarrow \varphi^* \circ \varphi_*v &= v \circ \varphi^*.\end{aligned}\tag{10.9}$$

So if v is an invariant vector field, we can write

$$\varphi^* \circ v = v \circ \varphi^*.\tag{10.10}$$

This expresses invariance under φ , and is satisfied by all differential operators invariant under φ .

Consider a vector field u , and the local flow (or one-parameter diffeomorphism group) ϕ_t corresponding to u ,

$$\phi_t(Q) = \gamma_Q(t), \quad \dot{\gamma}_Q(t) = u(\gamma_Q(t)). \quad (10.11)$$

But for any $f \in C^\infty(\mathcal{M})$,

$$\begin{aligned} \dot{\gamma}_Q(f) &= \frac{d}{dt} (f \circ \gamma_Q(t)) \\ &= \frac{d}{dt} (f \circ \phi_t(Q)) \\ &= \frac{d}{dt} (\phi_t^*(f)) = u_{\gamma_Q(t)}(f) \equiv u(f) \Big|_{\gamma_Q(t)} \end{aligned} \quad (10.12)$$

At $t = 0$ we get the equation

$$\frac{d}{dt} (\phi_t^*(f)) \Big|_{t=0} = u(f) \Big|_Q \quad (10.13)$$

We can also write

$$\frac{d}{dt} (\phi_t^* f)(Q) = u(f)(\phi_t(Q)) = \phi_t^* u(f)(Q). \quad (10.14)$$

This formula can be used to solve linear partial differential equations of the form

$$\frac{\partial}{\partial t} f(\mathbf{x}, t) = \sum_{i=1}^n v^i(\mathbf{x}) \frac{\partial}{\partial x^i} f(\mathbf{x}, t) \quad (10.15)$$

with initial condition $f(\mathbf{x}, 0) = g(\mathbf{x})$ and everything smooth. This is an equation on \mathbb{R}^{n+1} , so it can be on a chart for a manifold as well.

We can treat $v^i(\mathbf{x})$ as components of a vector field v . Then a solution to this equation is

$$\begin{aligned} f(\mathbf{x}, t) &= \phi_t^* g(\mathbf{x}) \\ &\equiv g(\phi_t(\mathbf{x})) \equiv g \circ \phi_t(\mathbf{x}), \end{aligned} \quad (10.16)$$

where ϕ_t is the flow of v .

Proof:

$$\frac{\partial}{\partial t} f(\mathbf{x}, t) = \frac{d}{dt} (\phi_t^* g) = v(f) \equiv v^i \frac{\partial f}{\partial x^i}, \quad (10.17)$$

using Eq. (10.13). \square

Thus the partial differential equation can be solved by finding the integral curves of v (the flow of v) and then by pushing (also called **dragging**) g along those curves. It can be shown, using well-known theorems about the uniqueness of solutions to first order partial differential equations, that this solution is also unique.

Example: Consider the equation in 2+1 dimensions

$$\frac{\partial}{\partial t} f(\mathbf{x}, t) = (x - y) \left(\frac{\partial f}{\partial x} - \frac{\partial f}{\partial y} \right) \quad (10.18)$$

with initial condition $f(\mathbf{x}, 0) = x^2 + y^2$. The corresponding vector field is $v(\mathbf{x}) = (x - y, -x + y)$. The integral curve passing through the point $P = (x_0, y_0)$ is given by the coordinates

$$\gamma(t) = (v_x(P)t + x_0, v_y(P)t + y_0), \quad (10.19)$$

so the integral curve passing through (x, y) in our example is given by

$$\begin{aligned} \gamma(t) &= ((x - y)t + x, (-x + y)t + y) \\ &= \Phi_t(x, y), \end{aligned} \quad (10.20)$$

the flow of v . So the solution is

$$\begin{aligned} f(\mathbf{x}, t) &= \Phi_t^* f(\mathbf{x}, 0) = f(\mathbf{x}, 0) \circ \Phi_t(x, y) \\ &= [(x - y)t + x]^2 + [(-x + y)t + y]^2 \\ &= (x - y)^2 t^2 + x^2 + 2(x - y)xt + (x - y)^2 t^2 + y^2 - 2(x - y)yt \\ &= 2(x - y)^2 t^2 + (x^2 + y^2)(1 + 2t) - 4xyt. \end{aligned} \quad (10.21)$$